



# RUNAWAY SPINS IN BINARY BLACK HOLES



## INTRO

- **Black holes** (BHs) are warps in spacetime so strong that they consume ALL things that get too close.
- In BH **binary** systems two BHs orbit each other.
- The emission of **gravitational waves** (GWs) leads the BHs to inspiral and eventually merge.
- The GW signal from a merging binary BH was first detected by LIGO in 2015:

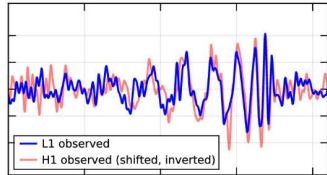


FIG 1: The signal GW150914, the first ever observation of GWs [2].

## PRECESSION

- Consider a spinning top. Unless spun perfectly upright, it wobbles in a cone as it rotates – this is called **spin precession**.
- BHs precess if the spins are not orthogonal to the orbit, significantly altering any GW signal.
- There are four of these aligned spin configurations:

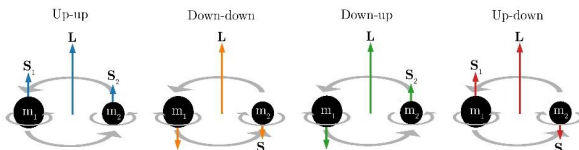


FIG 2: The four aligned spin configurations.  $L$  = orbital angular momentum,  $S$  = spin and  $m$  = mass.

- Are these systems **stable**? Or does spin precession change the configuration as the BHs evolve?

## THE MATHS

### 1 DEFINITIONS

- Mass ratio  $q = m_2/m_1$ , total mass  $M = m_1 + m_2$ , Kerr parameters  $\chi_1 = m_1^2|S_1|$ ,  $\chi_2 = m_2^2|S_2|$ .
- The spin vectors and their relative orientations are defined in the following figure and the evolution of binary BH spins is governed by the following equations:

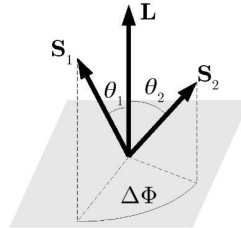


FIG 3: The relative orientations of the spin vectors  $S_1$  and  $S_2$  with the orbital angular momentum  $L$ , and the angles  $\theta_1$  and  $\theta_2$  between them.

$$\frac{dS_1}{dt} = \omega_1 \times S_1$$

$$\frac{dS_2}{dt} = \omega_2 \times S_2$$

$$\frac{dL}{dt} = \omega_L \times L + \frac{dL}{dt} \hat{L}$$

### 2 PERTURBATIONS

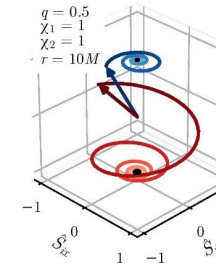
- To test the stability of the aligned spin configurations, we introduce small perturbations to the spin directions.
- This means the two spin vectors are slightly misaligned with the orbital angular momentum. We measure this misalignment with a parameter  $\epsilon$ .
- The perturbation evolves to leading order as a **simple harmonic oscillator**:

$$\frac{d^2\epsilon}{dt^2} + \omega^2\epsilon \approx 0$$

- The evolution of the perturbation depends on  $\omega^2$ :
  - $\omega^2 > 0 \Rightarrow$  the spins are **stable**
  - $\omega^2 < 0 \Rightarrow$  the spins are **unstable** and begin to precess
  - $\omega^2 = 0 \Rightarrow$  transition from **stability to instability**

## RESULTS

- From  $\omega^2$ , we find that only the **up-down** configuration (see FIG 2) is unstable to precession.
- Those binaries become unstable when the distance  $r$  between the BHs shrinks to values between:



$$r_+ > r > r_-$$

$$r_{\pm} = \frac{(\sqrt{\chi_1} \pm \sqrt{q\chi_2})^4}{(1-q)^2}$$

FIG 4: The evolution of the two BH spin vectors after an up-down binary encounters the instability.

- The **endpoint** of the unstable evolution is given by:

$$\cos\theta_1 = \cos\theta_2 = \frac{\chi_1 - q\chi_2}{\chi_1 + q\chi_2}$$

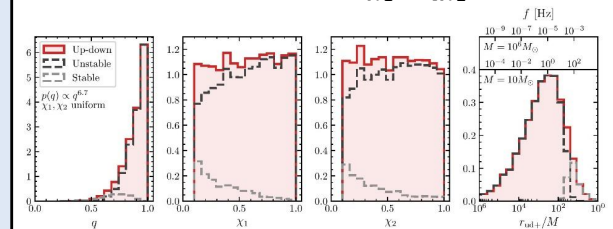


FIG 4: Astrophysical population of up-down BH binaries. Unstable binaries make up most of the population (left three panels) and undergo the instability before reaching the LIGO sensitivity window (frequency  $\geq 50$  Hz, right panel).

## REFERENCES

- [1] M. Mould and D. Gerosa, arXiv:2003.02281 [gr-qc] (2020).
- [2] B. P. Abbott *et al.* (LIGO Scientific Collaboration and Virgo Collaboration), Phys. Rev. Lett. 116, 061102 (2016).